

CHAPTER 9

Effects of Polymer Confinement in AAO Nanocavities on the Microstructure and Properties of Polymer Materials

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9.1. Introduction

Nanotechnologies require the combination of materials in which at least one dimension is nanoscopic, with specific chemical functions and easy nanofabrication processes. Among the range of materials available for nanotechnology applications, polymers are attracting a great deal of attention due to their customizable chemical and

1 physical properties and the fact that polymers can be easily processed
2 at the nanoscale. Hence, the field of polymer nanotechnology is cur-
3 rently experiencing a rapid development.^{1–6}

4 The replication of anodic aluminum oxide (AAO) templates is
5 establishing itself as a major approach for the fabrication of polymer
6 nanostructures, thanks to its large versatility, low cost, and high
7 yield. After the polymer material is subject to such a nanofabri-
8 cation process, the physical properties and behavior of the resulting
9 polymer nanostructures, namely, crystallization, molecular dynamics
10 and many others, are altered with respect to the properties of the
11 bulk polymer counterpart. This is due to the so-called “confinement
12 effects” which arise as a consequence of the low dimensionality of the
13 polymer material confined in the nanoporous AAO templates.^{7–11}
14 Thus, “confinement effects” become a new tool to modulate the
15 properties of polymer materials, thereby produce nanostructures
16 with tailored functions and morphologies, which are more and more
17 demanded for application as biosensors, scaffolds for cell adhesion,
18 hydrophobic surfaces, and others.²

19 In recent years, there have been hundreds of reports demonstrat-
20 ing the feasibility of AAO templates to produce polymer nanos-
21 tructures in a variety of polymers. For the fabrication of most of
22 these nanostructures, a pre-synthesized polymer is infiltrated inside
23 the pores either in form of melt or solution. A recently published
24 review article summarizes the strategies developed to produce poly-
25 mer nanostructures using AAO templates as well as a selection
26 of the most relevant morphologies achieved, such as nanofibers,
27 nanotubes and nanospheres, etc. made of polystyrene (PS), poly
28 (methyl methacrylate) (PMMA), poly (vinylidenefluoride) (PVDF),
29 and many other polymers and polymer-based composites.² Moreover,
30 that review gives detailed information concerning the influence of the
31 confinement on some structural and dynamical behavior of polymers,
32 such as the confirmation of amorphous chains, the glass transition
33 phenomenon and chain dynamics, as well as on the material’s prop-
34 erties — mechanical properties, transport properties, ferroelectric
35 properties, thermal properties, rheological properties, etc. A brief
36 section devoted to the use of AAO nanopores as nanoreactors for the

in situ polymerization, was also reported to suggest an alternative to the polymer infiltration of pre-synthesized polymers. We encourage the readers to have a look at that paper in order to gain further knowledge on these subjects.

The search for innovative preparation methods of novel polymer structures via sophisticated infiltration methods of polymers into nanopores is currently very active.^{12–17} In this chapter, we show the most recent results involving polymer patterning with AAO templates aimed to develop novel polymer nanostructures. Specifically, we provide updated information in Section 9.2.1 about the employment of AAO templates as nanomoulds for the preparation of two different hierarchically nanostructured polymers: core-shell nanocylinders by a double infiltration method and polymer nanotubes obtained by layer-by-layer (LbL) assembly of polyelectrolytes inside AAO membranes.

Moreover, we compile here the recent developments that are enabling fine control of the features of the confined polymer material at different length scales, namely at molecular level and at the mesoscale. In Section 9.3.1, we address the question of how the spatial restriction impacts on the molecular architecture of polymer chains synthesized by radical polymerization inside nanopores, with special focus on the molecular weight, Mw, and polydispersity of the synthesized polymer chains.

At a larger scale, i.e., at the mesoscale, the properties of polymer nanostructures are determined by their structural and compositional configuration. In Section 9.3.2, we review the recent advances on compositional design of polymer nanostructures produced via LbL deposition. Lastly, the impact of pore confinement on the structural development — solidification — of semicrystalline polymers is comprehensively described in Section 9.3.3.

9.2. Polymer Nanostructures Obtained from AAO Templating

Polymer nanostructures with modulated morphologies and properties can be easily achieved from the AAO templates assisted method.

This conventional method allows tailoring the dimensions of a huge number of polymer and polymer-based composites and, therefore, of their specific polymer properties and applications. In effect, the fact that polymer dimensions can be adjusted is a fundamental advantage in the study of size-dependency properties. Over the last years, a growing scientific interest in exploring new strategies to find new routes to fabricate polymer nanostructures have garnered great attention for specific purposes. In this section, we display three novel line of attack/approach to get in a straightforward and/or direct manner new architectures and morphologies: (i) new morphologies named core-shell nanocylinders, (ii) combination of different nano-micro methods, AAO-LbL process, and (iii) *in situ* nanosynthesized polymers.

9.2.1. AAO nanomolding

9.2.1.1. Double infiltration method (core-shell nanocylinders)

In the literature, have been reported different methods to prepare sophisticated polymer nanostructures from AAO templates.^{13,16} More recently, with the objective of modifying the surface of the alumina nanopores with a polymer coating to study polymer crystallization under no influence of AAO walls, a double infiltration method from the melt was developed.¹⁷ This method allows to fabricate by means of a solution-free method (greener and easy to implement), a series of core-shell nanostructures (nanocylinders) which can be the subject of different studies. In fact, this method provides an easier and more versatile way to study the crystallization or dynamics properties of a confined polymer, in comparison to other possibilities such as the chemical functionalization of AAO pores (silylation).

As an example of this methodology, a two-step polymer infiltration process was recently implemented to yield core-shell polycaprolactone-polystyrene (PCL-PS) nanocylinders. First, a PS nanotube inside the AAO template was generated by melt infiltration of PS in AAO templates of pore diameter of 350 and 140 nm, at around 200°C. Then, a second infiltration step of PCL was carried out to fill the nanotube and form the PCL core nanocylinder.

1 Nanotubes and filled nanotubes thickness and diameters can be
 2 determined by scanning electron microscopy (SEM) and transmis-
 3 sion electron microscopy (TEM). A correct infiltration of both poly-
 4 mers, meaning that the process resulted in no voids within the
 5 AAO templates was ascertained by fourier transform infrared spec-
 6 troscopy (FTIR) which also allowed to calculate the weight average of
 7 each polymer. SEM images corresponding to polymer nanostructures
 8 obtained from infiltration within AAO templates of 250 and 140 nm
 9 are shown in Figs. 9.1(a)–9.1(d) and 9.1(e)–9.1(f), respectively. In
 10 Fig. 9.1(a), PS nanotubes are observed; in Fig. 9.1(b), PCL-filled PS
 11 nanotubes; Fig. 9.1(c) corresponds to the core-shell nanocylinders
 12 (PS-shell and the PCL-core); Fig. 9.1(d) (the same image colored in
 13 orange (PS) and blue (PCL), as a guide the eye); and Figs. 9.1(e)
 14 and 9.1(f) correspond to PS nanotubes and PCL-filled PS nanotubes,
 15 respectively.

16 To summarize, PCL nanofibers are covered by a PS shell, so it is
 17 not under the chemical influence of AAO walls. The crystallization
 18 of PCL nanofibers in the bulk has been studied and compared with
 19 the crystallization of PCL nanofibers inside AAO templates of the
 20 same size, showing a different behavior. It was suggested a surface
 21 versus homogeneous nucleation.¹⁷

22 9.2.1.2. *LbL assembly (polymer nanotubes)*

23 The LbL procedure consists of the alternate adsorption of a polyanion
 24 and a polycation onto charged macroscopic surfaces giving rise to
 25 inversion of the superficial charge with every adsorption step so that a
 26 number of consecutive deposited layers give rise to polymer materials
 27 structured in layers.^{18–20} Over the past few years, there is growing
 28 interest in coupling the LbL methodology with the employment of
 29 porous substrates as templates for the preparation of nanostructured
 30 polymer materials with hierarchic functional structures.^{21,22} In this
 31 way, the precise pore geometry of the template is combined with the
 32 organized layered structure of an LbL assembly. Ordered templates
 33 include micro and nanofluidic devices with well-defined geometries,
 34 polycarbonate (PC) and AAO membranes.

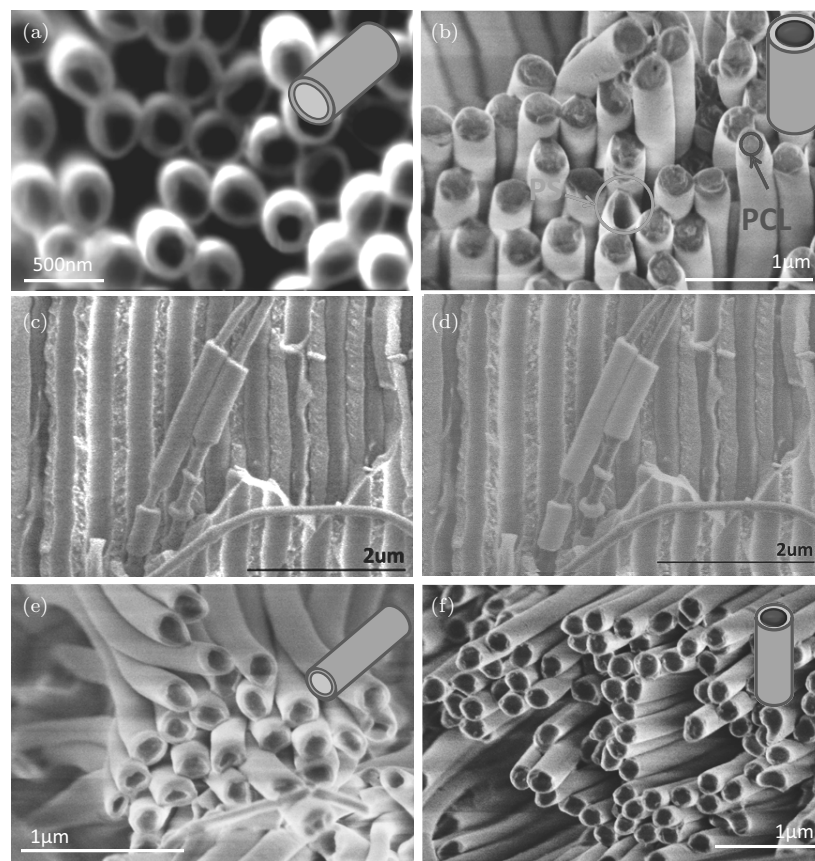


Fig. 9.1. SEM micrographs of (a) PS nanotubes after infiltration inside AAO templates with nanopores of 350 nm diameter, (b) PCL-PS core-shell nanostructure of 350 nm diameter (PCL core 220 nm diameter), (c) PCL-PS core-shell nanostructure of 350 nm diameter (PCL core 220 nm diameter) inside the AAO template (lateral view), (d) colored figure of (c) to guide the eye, blue for PCL core and orange for PS shell, (e) PS nanotubes after infiltration inside AAO templates with nanopores of 140 nm diameter, and (f) PCL-PS core-shell nanostructure of 140 nm diameter (PCL core 60 nm diameter).

Source: Reproduced with permission from Ref. [17].

- 1 A pioneering work from Caruso *et al.*, in 2003 showed for the first
- 2 time the possibility to combine the LbL assembly within the AAO
- 3 pores. LbL assembly on AAO membranes gives rise to cylindrical
- 4 nanotubes with controlled external diameters in which the length,

1 composition, and wall thickness can be modulated giving rise to
 2 multiwall nanotubes.²³ The process of LbL inside the nanopores of
 3 AAO membrane takes place experimentally by placing the membrane
 4 alternatively in solutions of the polycation and the polyanion as can
 5 be observed in Fig. 9.2.

6 In order to obtain free standing multiwall nanotubes, AAO mem-
 7 branes need to be dissolved in aqueous solutions at $\text{pH} < 4.5$ to
 8 $\text{pH} > 8.5$. These experimental conditions are pretty harsh toward

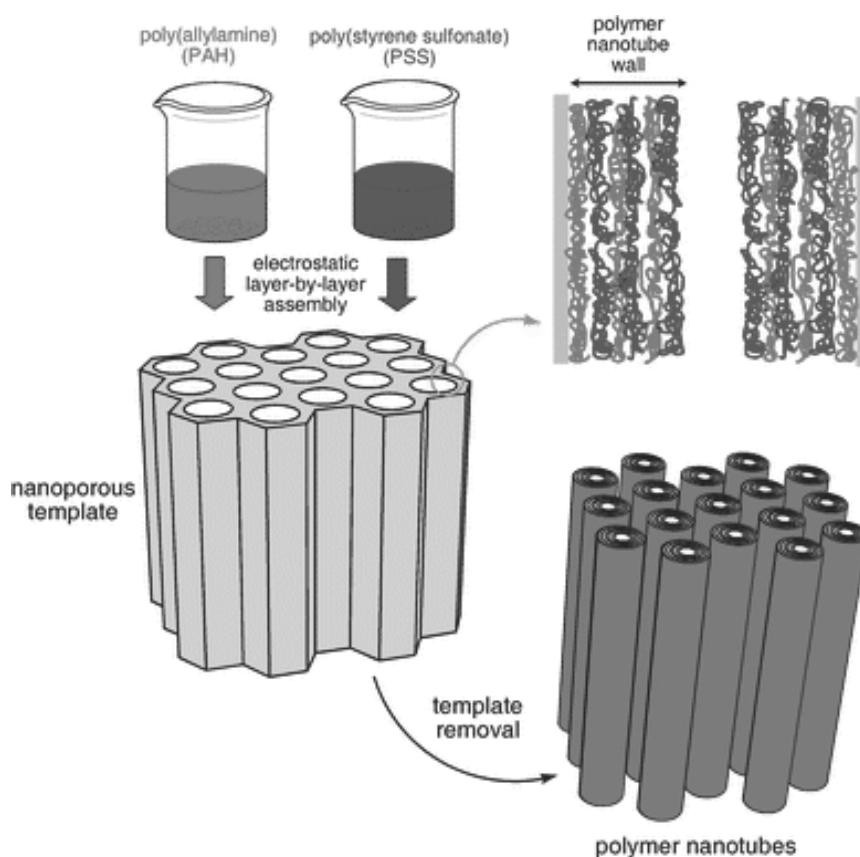


Fig. 9.2. Simplified representation of the formation of polyelectrolyte nanotubes through sequential assembly of polycations and polyanions inside the nanoporous template and the subsequent removal of the nanotemplate.

Source: Reproduced with permission from Ref. [22].

the polyelectrolyte assembly that holds together through ionic interactions. Alternatively, track-etched PC membranes which can be dissolved in mild solution conditions in organic solvents such as dichlorometane or dimethylformamide have also been employed for the study of the LbL assembly of different polyelectrolytes.²⁴

Still, there is a variety of materials that have been infiltrated within the cylindrical nanopores of AAO via the LbL assembly. One example is the deposition of protein-based nanotubes, whose chemical structure consists of polypeptidic chains with different ionizable groups for potential applications as biosensors and enzymatic bioreactors and for disease treatment and bioseparations.²⁵ As an example of preparation, AAO membranes were alternatively immersed in an aqueous solution of a protein followed by immersion in a glutaraldehyde solution, a crosslinking agent that facilitated the stability of the protein layers after AAO membrane dissolution. The number of protein layers that make up the nanotube walls can be tuned by changing the number of alternately deposited protein glutaraldehyde cycles.²⁶

TiO₂-based composite nanotube arrays with well-defined outer diameters and lengths, largely determined by those of the AAO templates, and wall thicknesses controlled by the number of layers deposited have been prepared for applications in catalysis, chemical sensors, nanoelectrodes, and nanodevices. To that aim, AAO membranes were functionalized by first adsorbing poly(sodium 4-styrenesulfonate), PSS, onto the surfaces to form the first layer, to render it negatively charged. After washing the substrate with water, positively charged poly(diallyldimethylammonium chloride), was infiltrated to give rise to a bilayer.²⁷

Many studies have focused on the fundamental behavior of polyelectrolyte deposition in confined geometries in order to determine differences with the process when carried out on standard, flat substrates.^{28–33} This topic will be developed in further detail, specifically focusing on LbL assembly on porous alumina templates, within Section 9.3.2.

9.2.2. AAO nanoreactor: *In situ* polymerization

Due to strong requirements imposed by polymer infiltration methods in some cases, high temperature and long time, the preparation of polymer nanostructures by *in situ* polymerization of a precursor monomer in AAO templates has recently emerged as an alternative and efficient route to the polymer infiltration method. In polymers, chemical structure is strongly dependent on the reaction environment since the highly reactive species, catalysts, radicals and anions, and their molecular diffusions involved in the polymerization process are strongly susceptible to the reaction environment. A clear example of this influence is in the biological systems, where natural polymers produced by enzymatic catalysis are developed within regulated and well-organized molecular scale spaces, such as, channel structures layered materials or nanolevel templates.³⁴ Therefore, the geometrical constraints of the well-defined nanochannels of a AAO template could affect the pattern of monomer insertion and chain growth process, thus offering a way to control the polymer chain structures and its macroscopic morphologies. As a consequence, of the great interest and importance on the subject, there are available in the literature quite a few studies on polymerization reactions in confined systems, such as, porous glass, porous coordination polymers, and the nanocavities of AAO templates, among others. Moreover, AAO nanochannels fulfill the requirements of a precision reactor, since the dimension can be adjusted and accuracy determined, therefore, this nanoreactor should be easily modeled.

In particular, the radical polymerization of vinyl monomers, such as, methyl methacrylate (MMA), styrene (St), and perfluoracrylate (FA) and the step-growth polymerization of a diol and a diisocyanate has been successfully achieved in the pores of AAO templates.^{35–39} Another illustrative example of polymerization in confinement has been reported by Tarnacka *et al.* upon step-growth polymerization of bisphenol-A diglycidyl ether (DGEBA) with aniline both in bulk and in AAO membranes.⁴⁰ In this case, the kinetic of reaction was followed by both FTIR and Broadband Dielectric Spectroscopies.

The first general observation is that kinetic curves do not follow sigmoidal shape that is characteristic for the autocatalytic type of chemical reactions. In a more detailed study, it was observed that (i) the polymerization was faster under confinement, compared to bulk reaction carried out at the same temperature conditions, (ii) the reaction speeds up with the degree of confinement, and finally, it was found that the initial step of the polymerization is significantly reduced or even suppressed in nanochannels. From the analysis of FTIR data with the support of Monte Carlo simulation they found that the rate of reaction is slower at the surface of the pore walls with respect to the polymerization at the core of nanochannels. Moreover, they found out that the activation barrier for the polymerization remains unchanged under confinement.

9.2.2.1. *In situ polymerization of MMA monomer in AAO templates*

The polymerization of this vinyl monomer is described in detail as an example of radical polymerization in confinement in AAO templates and the results compared to those obtained in the bulk. For the experiment, a solution of MMA and AIBN is placed on the top of a AAO template of 35 nm of pore diameter and 100 μm of pore length, then, the absorption of the solution into the template occurs quickly as a result of the capillary force and polarity of the solution. The reaction is carried out under isothermal condition at 80°C and monitored as a function of time by Raman microscopy, by following the decreasing of double bond bands, or by Differential scanning calorimeter (DSC) by measuring the area of exothermic and intense peak that appears associated to polymerization reaction. In Fig. 9.3(a) is plotted the polymerization kinetics of MMA at 60°C, 70°C, 80°C, and 90°C in bulk and within AAO of 35 nm pore diameter. As observed, the reaction within the AAO template is significantly faster compared to the bulk polymerization. It can also be observed that the initial rate of polymerization is much faster in the nanoconfined system than bulk and that the onset of the gel point occurs at slightly lower conversion. Moreover, the rate of polymerization after the gel point was

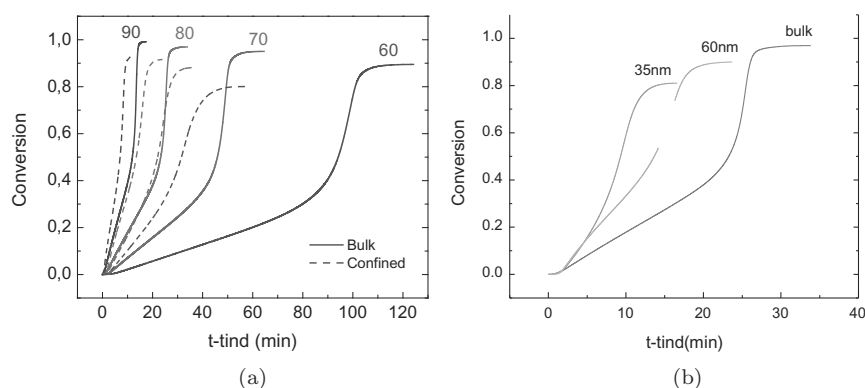


Fig. 9.3. (a) Kinetic plot of MMA polymerization in confinement at 60°C (blue), 70°C (turquoise), 80°C (pink), and 90°C (purple) (b) Evolution of conversion with time in the polymerization of MMA in bulk and within AAO templates of 35 and 60 nm.

1 slower and a lower limiting conversion is reached for polymerizations
 2 under nanoconfinement. The polymerization reaction of MMA has
 3 been studied within the AAO templates at different degrees of con-
 4 finements and results compared in Fig. 9.3(b). It is observed that as
 5 the degree on confinement increases (lower diameter) the reaction is
 6 faster. In an effort to parametrize the above differences, a mathemat-
 7 ical model was implemented and is explained in Section 9.3.

8 9.3. Confinement Effects on the Development 9 of Polymeric Architecture and Microstructure

10 Confinement effects of polymers within AAO templates present a
 11 strong influence on several intrinsic characteristics of polymers such
 12 as molecular weight or polydispersity and important physical prop-
 13 erties such as self-assembly, including crystallization. In this section,
 14 we provide important information on this important area of research.

15 9.3.1. Effects of confinement on the polymer 16 architecture

17 In the study of polymerization reactions in AAO templates, it is gen-
 18 erally observed that confinement effects strongly alter the reaction

1 kinetics and polymer structure, that is, molecular weight and molec-
 2 ular weight distribution. A general conclusion is that when a radical,
 3 controlled or step polymerization reaction is carried out in confined
 4 geometries, the polymerization kinetics and the resulting polymer
 5 structure are significantly altered.^{40–42} A mathematical model was
 6 implemented to study the free-radical polymerization of MMA in
 7 confinement in AAO templates and to explain the differences between
 8 the reaction kinetics and molecular weight for reactions carried out
 9 under nanoconfinement and bulk, in terms of confinement effects.
 10 Experimental data were fit using a mathematical model taking into
 11 account the following confinements effects: (i) Under confinement
 12 within the AAO walls, a faster decomposition rate of the initiator
 13 molecule leads to a faster rate of polymerization, (ii) Due to the
 14 higher glass transition temperature of polymers under nanoconfinement,
 15 the onset of diffusional limitations occurs earlier, and (iii) in
 16 the diffusion limiting region the pairwise combination of radicals also
 17 occurs faster in nanoconfinement than in bulk, due to the lower effective
 18 volume. Moreover, these confinement effects combine to give
 19 polymer of lower molecular weight and of lower dispersity.³⁷

20 The experimental molecular weight distributions (MWDs) of the
 21 chains can also be simulated, as shown in Figs. 9.4(a) and 9.4(b).
 22 The lower molecular weight region consists of chains that are formed

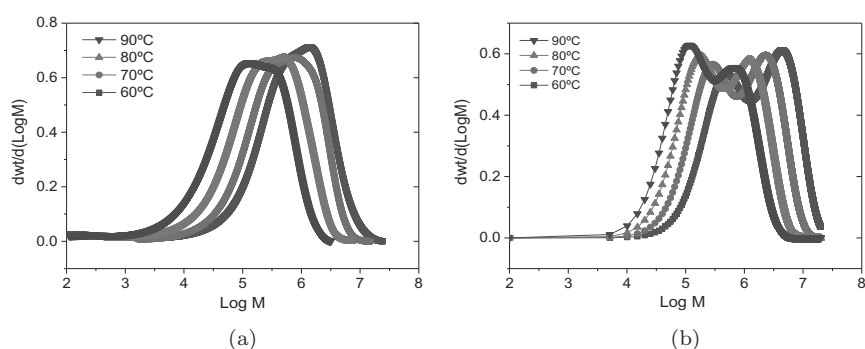


Fig. 9.4. (a) Experimental and (b) simulated molecular weight distributions of MMA polymerization in bulk at at 60°C (blue), 70°C (turquoise), 80°C (pink), and 90°C (purple).

1 in the early stages of the reaction, while the higher molecular chains
2 are generated during the second part of the reaction where, due to
3 the low rate of termination, the kinetic chain length is increased.
4 The simulated MWDs in Fig. 9.4(b) are in good agreement with the
5 experimental MWD, although due to broadening of the distribution
6 in the SEC measurement the bimodal nature of the experimental
7 distribution is poorly defined.

8 **9.3.2. Effects of confinement on polymer** 9 **self-assembly processes**

10 **9.3.2.1. Layer-by-layer assembly of polyelectrolytes**

11 The LbL assembly of polyelectrolytes within porous membranes can
12 be employed, on the one hand, as a technique for functionalization of
13 micro and nanoporous materials, including microfluidic devices^{43–45}
14 and, on the other hand for the formation of structured materials
15 employing the membranes as sacrificial templates.^{46,47} In any case,
16 there is a need to understand the mechanism of LbL in confinement
17 and how different experimental parameters as ionic strength, mem-
18 brane characteristics or polyelectrolyte concentration influence the
19 thickness of the resulting material.

20 AAO membranes provide the advantage of the control of the
21 pore size and thickness and thus they are the nanoporous membrane
22 of choice for fundamental studies regarding the effect of nanocon-
23 finement in LbL assembly of different polyelectrolytes. Still, the
24 complete understanding on the process of LbL growth on porous
25 substrates is lagging behind. This is mainly because of the diffi-
26 culty to carry out *in situ* measurements of film growth inside the
27 AAO nanopores. SEM is commonly employed to visualize infiltra-
28 tion of polymers inside AAO membranes, however, in the case of
29 infiltration of aqueous solutions of polyelectrolytes, the exact values
30 of the film thickness may not be measured from SEM images as
31 the vacuum environment in the SEM requires dried samples. There-
32 fore, there is a need to employ *in situ* experimental techniques to
33 monitor LbL polyelectrolyte growth within AAO membranes, mainly
34 optical techniques such as interferometric reflectance spectroscopy⁴⁸

or optical waveguide spectroscopy (OWS).⁴⁹ As an example, the *in situ* monitoring of the LbL deposition of polyelectrolyte dendrimers, dendritic structures with ionizable groups has been carried out by means of OWS.³² These materials constitute suitable model polyelectrolytes to follow the LbL assembly within nanopores because of their well-defined peripherally charged surface and their rigid internal hydrophobic structure that results in a lower degree of interlayer penetration than that encountered for linear polyelectrolytes.⁵⁰ Therefore, it is possible to determine the crossover point that defines the transition from normal assembly to assembly in nanoconfinement based on pore diameter.

Specifically, such study was carried out with fourth generation dendrimers (G4) having a theoretical diameter of 10 nm in a fully extended conformation and 7 nm in solution at high ionic strength under conditions of enhanced overall charge screening that allows an increase in the degree of dissociation. G4-polyelectrolyte LbL depositions were carried out on AAO membranes functionalized with a silane-derived molecule that provided positive surface charge. As can be observed in the SEM images corresponding to the interior and atop of a AAO membrane (65 nm) subjected to 10 deposition steps from 100 mM NaCl aqueous solution, the deposition takes place on the inside of the nanopore as well as on the surface of the AAO membrane (Figs. 9.5(a) and 9.5(b)). In addition, the pores in the AAO membrane were not completely filled. The thickness could be followed as a function of the number of dendrimer layers added separately for the top and the inner pore structure and the results are shown as a function of the AAO pore size in Figs. 9.5(c) and 9.5(d). The thickness of the film deposited on the surface of the AAO membrane increases linearly with the number of dendrimer layers added, similar to the trend found for the growth in planar surfaces. In addition, the growth rate is not affected by the presence of pores. However, in the case of the film infiltrated within the AAO nanopores, the film thickness linearly grows until a determined number of dendrimer layers after which a plateau in thickness is reached so that growth is inhibited beyond saturation values that decreased with pore size. That is, for 65 nm AAO membranes, thickness did not increase beyond

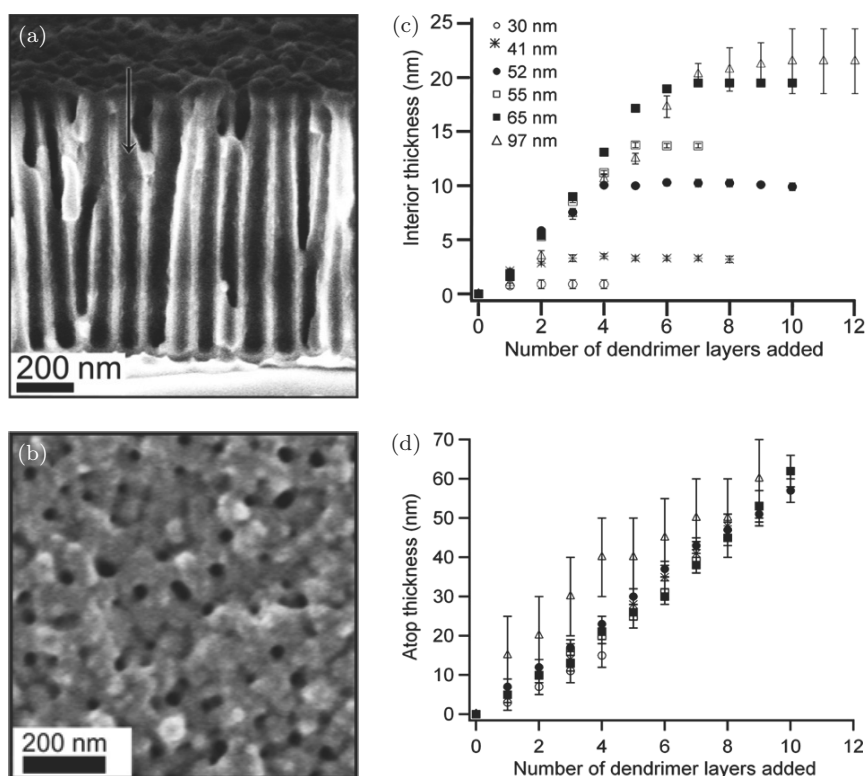


Fig. 9.5. SEM images of (a) the interior and (b) atop the AAO membrane ($D_0 = 65$ nm) after 10 deposition steps from 100 mM NaCl aqueous solution. (c) Thicknesses of the G4-polyelectrolyte layers deposited from 100 mM NaCl aqueous solutions, within (c) and (d) atop nanoporous AAO membranes with various pore diameters.

Source: Adapted from Ref. [32].

- 1 21.5 nm (6 LbL steps) while deposition within 30 nm pores immedi-
- 2 ately reached saturation during the first deposition step.
- 3 Rubner and coworkers studied the deposition of polymer/
- 4 nanoparticle composite multilayers within AAO membranes and
- 5 found that the presence of superficial charge in the AAO walls
- 6 gives rise to electrostatic repulsion along the pore diameter which
- 7 is enough to finish up the polyelectrolyte transport required for
- 8 LbL assembly.⁵¹ This would lead to the decrease of the level of the

1 compensation charge after every layer deposition, which would result
2 in lower thicknesses than that corresponding to the deposition on flat
3 substrates.

4 Other studies have demonstrated that the growth of multilay-
5 ers within the AAO membranes is independent from the molecular
6 weight or the ionic force of the solutions and it is only dependent on
7 the pore size at pore diameters < 250 nm.³⁰

8 Roy *et al.* demonstrated the existence of two growth regimes
9 for polymer nanostructures within nanopores.³³ The first regime is
10 similar to the growth found in planar surfaces. The polyelectrolyte
11 chains get close to the pore walls and get adsorbed similarly to the
12 deposition in planar surfaces. After the deposition of a number of
13 bilayers, the polymer chains start to interpenetrate inside the pores
14 giving rise to the second growth regime. The second regime is slower
15 in terms of kinetics due to the fact that diffusion is a limitant mech-
16 anism. The interpenetration of polyelectrolyte chains within AAO
17 pores might give rise to the formation of a dense gel that can clog
18 the pores and avoid more deposition steps. It is important to note
19 that the polyelectrolyte chains are highly hydrated and once they are
20 dried, the hydrated structures get shrunk resulting in a decrease of
21 the layer thickness.

22 In addition to a number of experimental studies about the LbL
23 assembly on porous substrates, there are few theoretical studies for
24 the molecular dynamic simulation of sequential adsorption of oppo-
25 sitely charged species carried out on charged substrates with cylindri-
26 cal pores of different sizes. In the case of modeling of the assembly of
27 nanoparticles, similarly to experimental results, the initial increase
28 in the film thickness is followed by a saturation regime. This was
29 attributed to the narrowing of the pore that inhibits penetration of
30 nanoparticles due to the electrostatic repulsion along the pore axis
31 that prevent the penetration of the nanoparticles deep inside the
32 pore.

33 As a final remark, a general conclusion that can be drawn from
34 different fundamental studies carried out on LbL assembly inside
35 AAO nanopores is that the polyelectrolyte deposition in confinement
36 is much more sensitive to the nature of the polyelectrolytes and the

1 ionic parameters compared to the deposition that takes place on
2 planar surfaces. In many cases, the solution ionic strength corre-
3 sponding to the infiltrated aqueous solutions is much greater than
4 that required for experiments on flat surfaces in order to efficiently
5 infiltrate the polyelectrolytes within the nanopores. In addition, the
6 final thickness resulting from the assembly of polyelectrolytes is sig-
7 nificantly lower in confined geometries than in standard, flat sub-
8 strates.

9 9.3.2.2. *Polymer crystallization*

10 Crystallization can be considered as the most common self-assembly
11 process in nature.⁵² Polymers typically crystallize as lamellar crys-
12 tals in which individual chains assemble normal to the surfaces of
13 lamellae adopting helical conformations.^{53,54} The thickness of lamel-
14 lae is typically around 10 nm⁵⁵ and develops approximately in the
15 polymer chain direction. Along the in-plane directions, in contrast,
16 the lamellae can grow up to micrometers (even millimeters). On a
17 next level, crystalline lamellae commonly pile up separated by disor-
18 dered amorphous regions,⁵⁶ which are comprised of entangled chain
19 segments, chain-ends, chain branches, chemical defects, etc. The set
20 of lamellae tend, in turn, to organize in larger polycrystalline assem-
21 blies; for example, densely branched spherulites with sizes that can
22 exceed millimeters^{52,55} are generally achieved when bulk semicrys-
23 talline polymers are cooled down from the melt.

24 While the thickness of lamellae is comparable to the pore diam-
25 eter of AAO templates, their lateral dimensions — and thus, the
26 size of the spherulites — surpass by far the pore diameter values.
27 Therefore, when a polymer melt is confined in AAO nanopores, its
28 natural crystallization process is altered. As a matter of fact, the
29 lateral size of the pores (defined by the diameter) conflicts with the
30 characteristic length scales of both processes the crystallization is
31 composed of, i.e., the crystal nucleation process and the crystal-grow
32 process, which result in drastic changes in the crystallization kinet-
33 ics and the microstructure (e.g., texture,) of the confined polymer as
34 compared to the same material with no space limitation.

(a) Crystallization Kinetics

Crystallization kinetic studies are typically accomplished monitoring the isothermal solidification of the polymer melt, such as a thermal profile yields a material with a more uniform inner structure than that of polymers crystallized under a constant cooling rate. Moreover, the evolution of the degree of crystallinity with time under isothermal conditions can be modeled by relatively simple mathematical models, the Avrami model⁵⁵ being the most frequently employed one. In the Avrami model, the advance of the crystallinity with time, i.e., the crystallization kinetics, can be expressed as:

$$V_c(t) = 1 - \exp(-Kt^n), \quad (9.1)$$

where $V_c(t)$ is the crystallinity in terms of the volume fraction of crystals (or supracrystalline structures) that develop within a specific time t ; K is a rate constant of the overall crystallization process, and n is the Avrami exponent. It is clear from Eq. (9.1) that the parameters n and K govern the crystallization kinetics.

(b) The Avrami exponent, n

The Avrami exponent, n , relates to the time dependency of the crystallization process; thereby, provides insights about every aspect of the crystallization process that depends on time. Consequently, n can be, in turn, interpreted as

$$n = n_n + n_g, \quad (9.2)$$

where n_n and n_g are the contributions to n resulting from how the nucleation process and crystal growth process depend on time, respectively. Typically, n_n ranges between zero to one. For example, in purely instantaneous nucleation processes where the large majority of nuclei develop at the same time, such as in the athermal heterogeneous nucleation occurring in bulk polymers, $n_n = 0$. Conversely, $n_n = 1$ when nuclei are developed at constant rate, like it is frequently assumed in purely homogenous nucleation processes.

On the other hand, n_g informs about the dimensionality of the growth of crystals when a constant growth rate is assumed, and adopts the values one, two or three when crystals (or supracrystalline structures) develop in one, two or three spatial dimensions, respectively.

Bulk polymers typically crystallize via heterogeneous nucleation mechanism ($n_n = 0$) followed by a 2D or 3D spherulitic growth of crystals ($n_g = 2$ or 3).⁵⁵ Thus, n ($n = n_n + n_g$) acquires values of two or three (or even higher) that result in the typical sigmoidal curves for the advance of the crystallinity ($V_c(t)$) with time that are commonly observed in bulk polymer crystallization.⁵⁵ In contrast, the value of n drops in crystallizations developed within pores as will be described in what follows giving rise to changes of the functional form of $V_c(t)$.^{57–60}

The main reason for the reduction of the value of n in crystallizations within pores is that n_g decreases strongly because of the restricted volume available for crystals to grow. Note that the maximum theoretical dimensionality for the growth of crystals in (highly confining) cylindrical pores is one, which is just achieved when the crystals can freely propagate straight on the pore.^{61,62} When the growth of crystals along the pore is impeded, $n_g < 1$ can even be achieved. The reduction of n_g is in general large enough to overcome the typical increase of n_n induced by the transition from heterogeneous nucleation mechanism in bulky polymers ($n_n = 0$) to homogeneous nucleation mechanism that is frequently active in confined systems (where $n_n = 1$, as will be described in detail below).^{59,63–69} Moreover, the dominance of the nucleation process over the crystal growth process in crystallization of confined polymers also contributes to lowering n , as n_n dominates against n_g , and n_n ranges solely between zero and one.

Michell *et al.* studied the crystallization of poly(ethylene oxide) (PEO) inside AAO pores of 60 and 35 nm in diameter. As shown in Fig. 9.6(a), n in these crystallizations were of close to one (between 0.7 and 1.4),⁵⁷ which was rationalized in terms of a complete suppression of the crystal growth due to the limited space ($n_g = 0$)

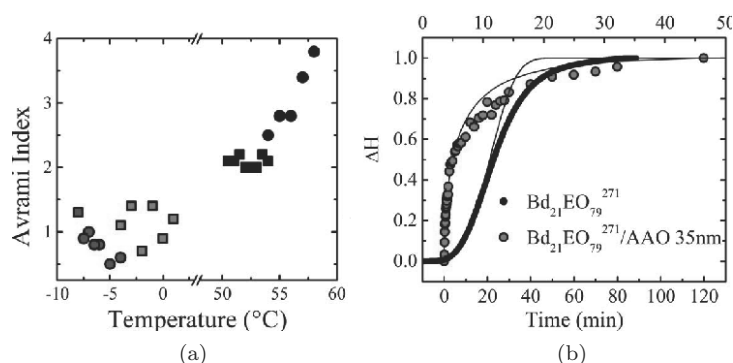


Fig. 9.6. (a) Avrami index versus T_c for bulk PEO (black squares), PEO confined in pores of 60 nm (red squares), PEO confined in pores of 35 nm (blue squares), bulk P(Bd-b-EO) copolymer (black circles), and P(Bd-b-EO) copolymer confined in pores of 35 nm (blue circles). (b) Advance of crystallinity (expressed as the relative melting enthalpy values, ΔH) with time for bulk PEO-b-PB (black circles, crystallization temperature (T_c) was 58°C), and within 35 nm AAO templates ($T_c = -4^\circ\text{C}$). The solid lines correspond to fits to the Avrami equation. The top x -axis corresponds to the bulk materials while the bottom x -axis corresponds to that of the confined material.

Source: Extracted from Ref. [58].

1 and homogeneous nucleation mechanism ($n_n = 1$). In further stud-
 2 ies, even lower n values (0.5–1) were reported by Michell *et al.* for
 3 PEO domains in a poly(ethylene oxide-*b*-butadiene) (PEO-*b*-PB)
 4 diblock copolymer⁵⁹ (Fig. 9.6(a)). These crystallizations exhibited
 5 nearly-first-order-kinetics and, subsequently, an exponential trend
 6 for the crystallinity with time (Fig. 9.6(b)). Similar arguments were
 7 employed by Maiz *et al.* to explain the reduction of the n values
 8 (ranging from 1.3 to 2.1) in the crystallization kinetics of polyethy-
 9 lene (PE) confined in AAO nanopores.⁶⁴

10 Note that in all these studies, n_g is the main contribution to
 11 n and, unlike in typical crystallizations of bulk polymers, acquires
 12 a non-zero value. The rationale for the non-zero value of n_g is
 13 that, under confinement, a transition frequently occurs from instan-
 14 taneous, heterogeneous nucleation of neat bulk polymers to a
 15 heterogeneity-free, homogeneous nucleation process, which is time-
 16 dependent.^{57–59,61,63–77} Bulk polymers typically crystallize via a

heterogeneous nucleation mechanism originated on heterogeneities (impurities) that are present in the melt. The typical density of impurities able to cause the nucleation of common thermoplastics, such as polypropylene (PP) and poly(vinyl difluoride) (PVDF), lies in the range of 10^6 – 10^7 cm⁻³.⁷⁸ This means that, in average, there is one heterogeneity every 10^5 – 10^6 μm³ of crystallizable material. However, the volume of a crystallizing polymer domain in a pore with a diameter of 35 nm and a depth of 100 μm is of $\sim 10^{-1}$ μm³. Thus, the number of crystallizing domains containing heterogeneities is insignificant and, hence, bulk-like heterogeneous nucleation mechanism is residual in these samples. The confined polymer nucleates, thus, through a different, heterogeneity-free mechanism, such as the homogeneous nucleation mechanism.

Interestingly, Maiz *et al.* found also first-order crystallization kinetics ($n = 1.1$ and 1.2) in the isothermal crystallization of PEO nanotubes with an outer diameter of 400 nm and a tube-wall of 35 nm;⁷¹ however, they observed also a significantly lower value of the apparent activation energy of the crystallization of the nanotubes compared to that of the bulk PEO, which suggests that their nanotubes were not crystallizing from homogeneous nucleations. Therefore, the first-order-kinetics in this system was interpreted as arising from the heterogeneous nucleation of crystals at the pore walls of the AAO templates, followed by the 1D growth of crystals.

Lastly, it is worth mentioning that pores exceeding 100 nm in diameter seem to have little impact on the crystallization kinetics of polymer melts.^{63,76,77}

(c) The crystallization rate constant, K

K being the rate constant for the overall crystallization process (see Eq. (9.1)), it is influenced by the nucleation rate, the number of nuclei, and the crystal growth rate. Among these factors, the major differences between confined and non-confined crystallizations are found in the number of nuclei. If, as it has been described above, the number of heterogeneities in bulk polymers is typically of 10^6 – 10^7 cm⁻³, in average, a crystalline nucleus is developed every 10^5 – 10^6 μm³. However, the volume of the crystallizable domains, i.e.,

that of the nanopores, lies in the range of $\sim 10^{-1} \mu\text{m}^3$. Thus, even in the unlikely case in which just a single nucleation center is developed inside each pore, the total density of nuclei in the polymer/AAO ensemble would be six to seven orders of magnitude higher than in typical bulk polymers (the density of pore in AAO templates is $10^8\text{--}10^{11}$ pores/ cm^{-2} ,^{79–82}). This dramatic increase of the number of nuclei in polymers confined in AAO templates provokes a substantial increase of K in these systems, as compared to bulk polymers (at least at low supercoolings). Note that, like n , K is also governed by the nucleation stage in crystallizations in confinement.

Woo *et al.*⁶⁹ reported for the first time that the K values of PE confined in AAO pores were five orders of magnitude higher than those of unconfined, bulk PE. Likewise, Michell *et al.* found an increase of two orders of magnitude in the K values of PEO confined in 60 nm pores with respect to bulk PEO.⁵⁷

Because the values of K are largely governed by the early stages of the crystallization process in confined polymers, K and the half-crystallization time ($t_{1/2}$) exhibit similar trends when plotted against the isothermal crystallization temperature, T_c (once the values of K are elevated to the power n^{-1}). Hence, $t_{1/2}$ can be also understood as a measure of the isothermal crystallization rate in confined systems and compared, thus, to K . As a matter of fact, Michell *et al.* showed the corresponding increase of the values of $t_{1/2}$ for the same PEO system confined in 60 nm pores with respect to its bulk counterpart.⁵⁸ Likewise, Maiz *et al.* found higher $t_{1/2}$ values in the crystallization of PE inside pores of 35 nm in diameter than in bulk PE.⁶⁴

(d) Crystal orientation (texture)

One of the most striking effects of pore-confinement is that high degrees of crystal orientation are induced. Furthermore, in many cases, the texture of the achieved nanostructure can be selected via suitable selection of the pore size^{72,83} and the presence or absence of a bulk polymer reservoir^{68,84,85} at the surface of the AAO template. It should be pointed out, though, that here we mainly focus on polymer nanostructures crystallized in the absence of such bulk reservoir.

The general trend in the crystallization of polymers confined in the straight cylindrical nanopores of AAO templates is that crystals propagate laying the polymer chains normal to the long axis of the nanopores.^{61,64,65,73,75,77,83,84} The rationale for this orientation is that polymeric chain-folded-lamellae can just grow along lateral directions, i.e., directions that are perpendicular to chains. Hence, the orientation of chain normal to the long axis of the pores allows the lamellae to propagate straight on the pores. Using Miller index notation, the crystals having $\langle hkl \rangle$ directions with a zero l index are the ones growing preferentially along the pore axis direction, as the l index (i.e., the c -axis of the crystal lattice) is approximately parallel to the chain direction in most of the polymer unit cells. Note, therefore, that, whereas the nucleation stage dominates the crystallization kinetics, the texture tends to be governed by the kinetic aspects of the crystal grow process.

A clear example of this texture has been recently found by Martín *et al.* in PVDF nanotubes.⁸⁶ Figure 9.7 shows the WAXS patterns of bulk PVDF and PVDF nanotubes still embedded in the AAO template (Fig. 9.7(a)). In these measurements, the scattering vector, q , was parallel to nanopore long axis. In the 2θ -range investigated, the bulk PVDF pattern showed all the characteristic reflections of α -PVDF, namely (100), (020), (110), and (021),⁸⁷ which indicates that the crystalline moieties in these samples are isotropically oriented. The pattern of the PVDF nanotubes also contains only features for α -PVDF, however, it displayed only reflections from the (020) and (110) lattice planes. Thus, crystals with $\langle hkl \rangle$ directions with a zero l -index are solely detected along the nanotube long axis direction.

In order to gain further information about the texture of PVDF nanotubes, XRD pole figure measurements were accomplished for the (020), (110), and (021) reflections (Fig. 9.7(b)). In these, the ψ -angle, defined as the angle produced by the rotation of the sample around an axis parallel to the plane of the template surface, is plotted along the radial direction, while the ϕ -angle, defined as the angle produced by the rotation of the sample around the long axis of the nanotubes, is plotted along the azimuthal direction (see inset in Fig. 9.7(a) for the identification of the angles). The (020) reflection centered at $\psi = 0^\circ$

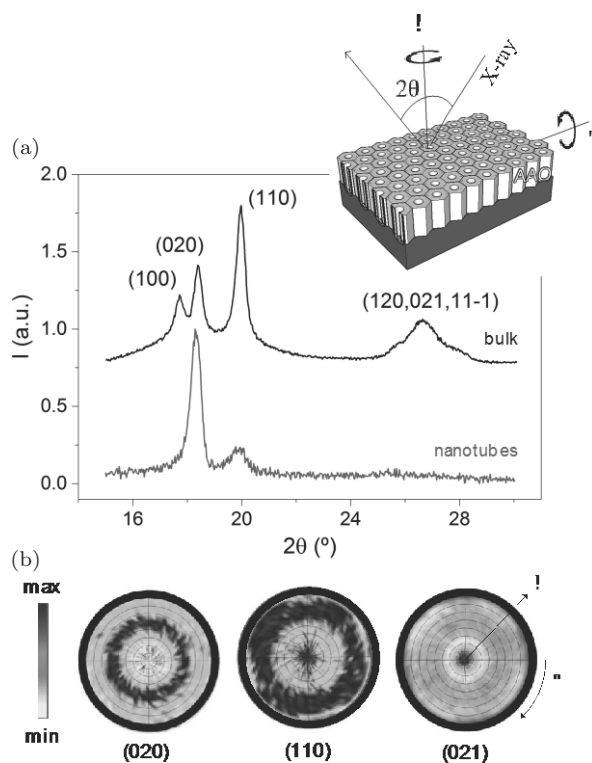


Fig. 9.7. (a) WAXS patterns of bulk PVDF and PVDF nanotubes, when q is parallel to the long axis of the nanotubes, i.e., $\psi = 0^\circ$. The geometry of the experiments as well as the relevant, scanned angles are identified in the schematic (inset). (b) XRD pole figures of the (020), (110), and (021) reflections measured in PVDF nanotubes, where the radial direction corresponds to the ψ -angle and the azimuth corresponds to ϕ .

Source: Adapted from Ref. [86].

revealed that the nanotubes were mainly comprised of crystals in which the crystallographic $\langle 020 \rangle$ direction was aligned with the nanotube long axis. These crystalline lamellae are isotropically oriented along the azimuthal ϕ -angle, as deduced from the ring-like patterns of (110) and (210) reflections.

Similar results were obtained by Martín *et al.* for the semiconducting polymer poly(9,9-dioctylfluorene) (PFO) confined into 35 nm AAO pores.⁸⁴ The 2D-WAXS pattern of this sample was

characterized by well-defined reflection at $q = 1.48 \text{ \AA}^{-1}$ (d -spacing of 0.425 nm) that is ascribed to the stacking of (530) lattice planes of the α -phase of PFO. A less intense reflection, corresponding to the stacking of (200) lattice planes, was also observed at $q = 0.49 \text{ \AA}^{-1}$. Thus, most of crystals pointed their (530) crystallographic direction parallel to the pore axis, although a small crystal fraction existed, where the $\langle 200 \rangle$ direction was aligned to the pore axis. Note that, again, both crystallographic directions corresponded to the $\langle hkl \rangle$ directions with zero l -index and as such, the chains lay normal to the longitudinal axis of pores.

Highly texturized polymer nanostructures with a single, uniaxial orientation can be also achieved when polymer melts are crystallized inside AAO porous templates. The formation of these ultra-oriented structures seems to be associated with a low nucleation rate as well as the existence of a crystallographic direction having a significantly faster growth rate than rest of the directions. Thus, under low nucleation rate conditions, such as low cooling rates in dynamic crystallizations or high T_c in isothermal crystallizations, some polymers crystallize aligning just the crystallographic direction with the fastest growth rate parallel to the longitudinal axis of the pores. As proposed by Huber,^{88–90} this can be traced to a crystallization mechanism, first suggested by Bridgman, for the single crystal growth in narrow capillaries.⁹¹ Within small capillaries, the crystalline direction with higher growth rate propagate along the direction of the long axis, while other growth directions die out.

By means of 2D-WAXS, Michell *et al.* investigated the texture of PEO confined in pores of 20 nm in diameter.⁶⁶ Figure 9.8 shows the 2D-WAXS patterns obtained when the X-rays traveled along the x -axis (a), y -axis (b), and z -axis (c) of the sample, the z -axis being parallel to the long axis of the nanopores, and x and y two directions mutually orthogonal and normal to the z -axis. The 2D-WAXS patterns resulting from the scattering of the PEO nanostructures in the z - x and the z - y planes (Fig. 9.8(a) and 9.8(b), respectively) showed the diffractions ascribed to the (120) planes located mainly along the meridian. Conversely, the 2D-WAXS pattern corresponding to the x - y plane displayed ring-like signals for all the diffractions.

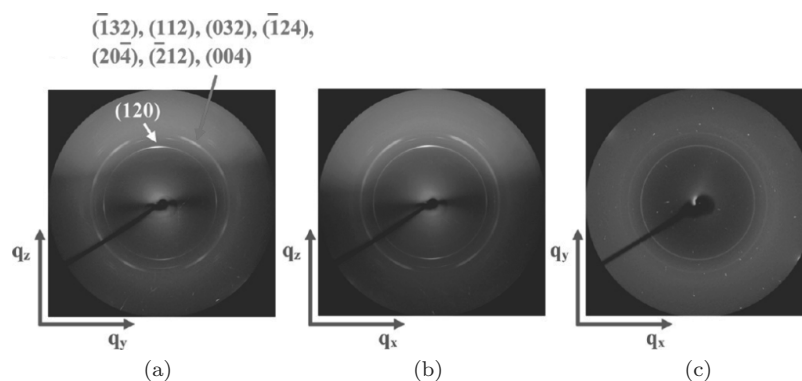


Fig. 9.8. 2D-WAXS patterns of PEO confined in 20 nm pores, collected for the X-ray beam traveling along the x -axis (a), y -axis (b), and z -axis (c) of the sample, the z -axis being parallel to the long axis of the nanopores, and x and y two directions mutually orthogonal and normal to the z -axis.

On the basis of this results, they concluded that PEO crystal moieties laid their $\langle 120 \rangle$ direction mainly parallel to the longitudinal axis of the pores but they were randomly oriented with respect to the x - y plane.

Perfectly uniaxial PEO nanorods were achieved by Martín *et al.* employing a crystallization approach based on inducing heterogeneous crystal nuclei at one of the tips of the nanorods.⁶¹ By the inducement of efficient surface interactions between the nanoconfined polymer and a material with nucleating capacity at the pore mouths (metallic Al, see SEM images in Fig. 9.9(a) and 9.9(b)), they managed to generate heterogeneous crystal nuclei at temperatures much higher than natural nucleation temperature of the confined polymer, which was strongly depressed because of the homogeneous nucleation mechanism. Thus, when the samples were isothermally crystallized at temperature slightly below the new nucleation temperature, no nucleation centers could develop in the confined material other than those at the pore mouths. Hence, the confined PEO nucleated solely at one of the extremes of the nanopores and then the crystallites propagated unidirectionally along the pores.

These PEO nanorods exhibited perfectly uniaxial crystal orientation, where the $\langle 120 \rangle$ lattice plains of the monoclinic cell unit of

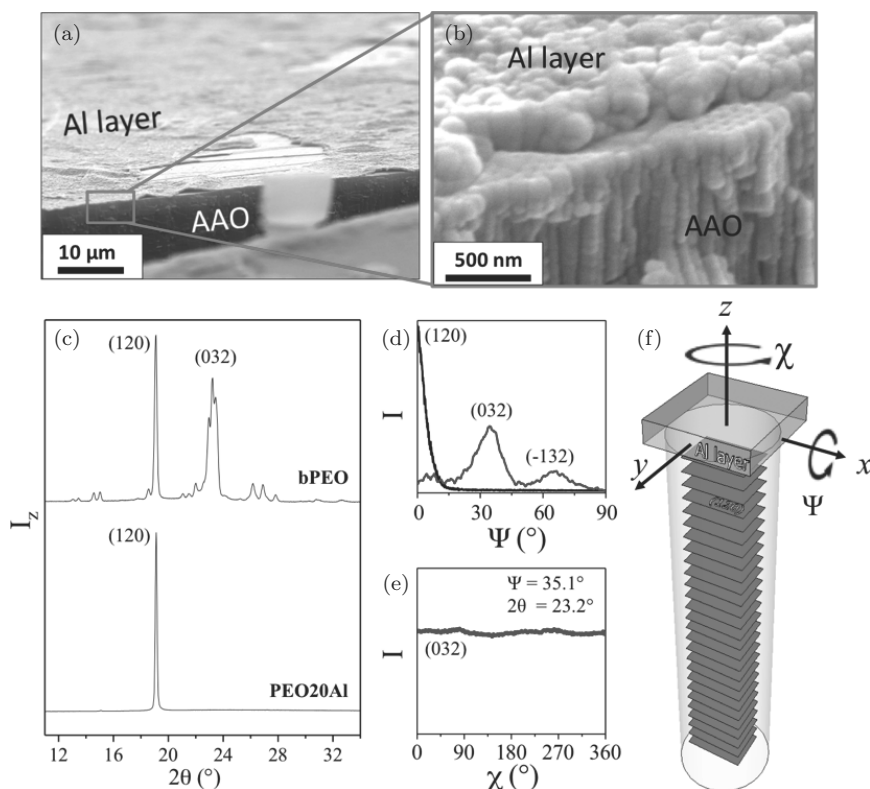


Fig. 9.9. Large view (a) and detailed view (b) SEM micrographs showing the Al layer deposited onto the AAO. (c) WAXS $\theta/2\theta$ scan of bulk PEO (bPEO) and confined PEO (PEO20Al) analyzed in the geometry where q was parallel to the PEO nanorods (z direction). (d) WAXS intensity recorded at $2\theta = 19.1^\circ$ (black line) and $2\theta = 23.2^\circ$ (blue line) while tilting the sample around the axis defined by the intersection of template surface plane and X-ray plane of incidence (Ψ angle). (e) WAXS intensity recorded along the angle confined within the plane of the template surface for the (032) reflection. 2θ was fixed to 23.2° , and Ψ was fixed to 35.1° . (f) Schematics of the crystal orientation PEO crystals in the pores (red foils represent the (120) planes of crystal lattice).

Source: Adapted from Ref. [61].

- 1 PEO were normal to pore axis (Fig. 9.9(c)). The orientation dis-
- 2 tribution of (120) planes was studied collecting WAXS intensity as
- 3 tilting the sample along the Ψ -angle (see the schematic in Fig. 9.9(f))
- 4 while setting the 2θ angle to the position of the (120) reflection. The

WAXS signal thus acquired presented a sharp maximum centered at $\Psi = 0^\circ$, indicating that the corresponding lattice planes lay exclusively perpendicular to the pore axis as illustrated in Fig. 9.9(f). This was further evidenced by analyzing the orientation distribution of the (032) and (−132) diffractions, which exhibited maxima just at the theoretical angles between the respective lattice planes and the (120) lattice plane (Fig. 9.9(d)).

Note that all the highly texturized structures presented here have in common that the nucleation of crystals within the pores was effectively reduced. However, there is an extreme case where the nucleation within the pores is even suppressed. This corresponds to a crystallization scenario where the polymer melt contained into the nanopores is connected through a bulk surface reservoir of material, so that the polymer contained in each pore are no longer isolated entities. In such a case, the volume of the total crystallizing material is by far larger than the average volume required for the presence of a heterogeneity. Thus, upon cooling from the melt, this polymer system behaves as a bulky system, crystallizing at the typical bulk temperatures via heterogeneous nucleation. Since the probability of finding heterogeneities in the surface reservoir is significantly higher than in the nanopores due to the larger volume of the former, the heterogeneous nuclei developed first in the reservoir and then grow isotropically through it. Some of these crystals grow in the direction of the pores and are thus able to propagate through the material contained in them. Therefore, the confined polymer crystallizes by the growth of crystals nucleated outside of the pores, i.e., no further nucleation events develop within the pores. The texture of the polymer nanostructures thus crystallized is governed solely by the crystal growth process, in such a way that, again, the “Bridgman-mechanism”⁹¹ for the propagation of the fast-growing crystallization mode along the pore determines the crystal orientation. Hence, perfect uniaxial orientations are frequently achieved when this crystallization strategy is used.^{64, 68, 92, 93}

The possibility to select crystal orientations by using surface reservoirs was first discovered by Steinhart *et al.*⁶⁸ Further on, Maiz *et al.* observed the same effect for PE crystallized with surface

1 reservoir inside 60 nm in diameter AAO templates.⁶⁴ Their WAXS
2 analysis revealed a pronounced alignment of the $\langle 020 \rangle$ crystallo-
3 graphic direction with the long axis of the pores. Interestingly, the
4 same publication demonstrated that when PE is forced to crystallize
5 in the absence of surface reservoir, $\langle 020 \rangle$, $\langle 110 \rangle$, and $\langle 200 \rangle$ crystal-
6 lographic directions align with the long axis of the nanopores. Note
7 that, in these configurations the chains are laid perpendicular to the
8 long axis of the nanopores, but do not correspond to a uniaxial crystal
9 orientation.

10 Uniaxial orientations were also found by Martín *et al.* for poly(3-
11 hexylthiophene) (P3HT) inside 250 nm and 350 nm pores. These
12 P3HT nanostructures crystallized orienting the b -axis of the orthor-
13 rhombic P3HT crystal parallel to the long axis of nanopores.^{83,94} The
14 b axis is the direction along which the π - π stacking occurs and,
15 thus, it is well known to be a fastest growth direction of the crystal.
16 Interestingly, they found that the presence of crystals with the $\langle 100 \rangle$
17 direction parallel to pore axis increased with the reduction of the
18 pore diameter. They hypothesized that the semirigid polymer chains
19 adopt flat, face-on conformations at the interface with the pore walls,
20 which would be responsible for the $\langle 100 \rangle$ crystal orientation.

21 In line with this result, further recent works have suggested that
22 strong confinement conditions^{72,83} (i.e., extremely narrow pores),
23 high chain rigidity,⁸⁴ and surface nucleation processes⁷¹ may induce
24 crystal textures other than those with the c -axis of the crystal ori-
25 ented normal to the pore long axis. Thus, Guan *et al.* reported that
26 inside pore with diameters above 10 nm, PEO crystals were oriented
27 with the $\langle 120 \rangle$ direction parallel to pores, whereas within 10 nm diam-
28 eter pores crystallites grew orienting chains parallel to the long axis
29 of the pores.⁷² They rationalized this phenomenon assuming that
30 when the pore diameter is smaller than the chain's contour length,
31 molten chains adopt anisotropic conformations to accommodate to
32 the reduced space, which upon crystallization, template the orienta-
33 tion of crystals.

34 Similar results were reported by Maiz *et al.* for PEO nanotubes
35 having 400 nm outer diameter and a tube-wall of 35 nm.⁷¹ Their
36 WAXS analysis revealed that crystals with $\langle hkl \rangle$ directions with

1 non-zero l -index where majority along the nanopore main axis. As
2 mentioned in the previous section, they also found signs of het-
3 erogeneous nucleation in these nanotubes; thereby, they proposed
4 a crystallization scenario where the nuclei were generated at the
5 PEO-AAO interface and then crystals grew mainly in the radial
6 direction of nanotubes until they reach the inner wall, where they die.

7 **9.4. Concluding Remarks**

8 Polymer 1D nanostructures derived from AAO templates, such as
9 polymer nanowires and nanotubes, have demonstrated great poten-
10 tial as building blocks to produce hierarchically organized nanos-
11 tructures. The recent research developments in the field are allowing
12 fine control of the confined polymer features at the different length
13 scales relevant for materials properties; i.e., at molecular level via
14 *in situ* polymerization and, at the mesoscale, through crystallization
15 engineering and compositional design via LbL deposition and multi-
16 component infiltration. Hence, further advances are expected in the
17 future that will allow customization of properties in the 1D nanos-
18 tructures as well as their integration into nanostructured devices that
19 may lay the foundation of a future technology with applications in
20 sensing, electronics, photonics, biomedicine, etc.

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